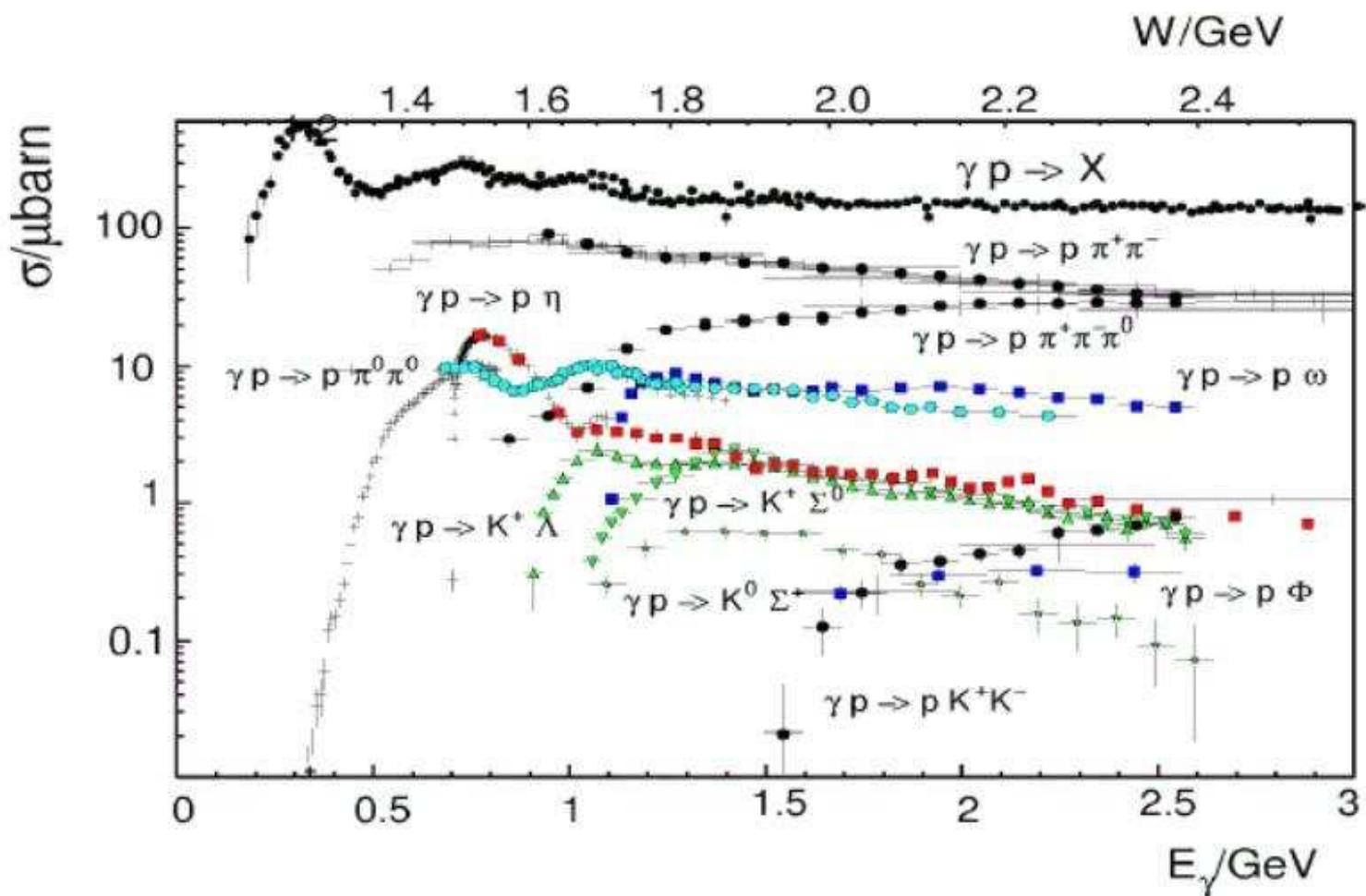
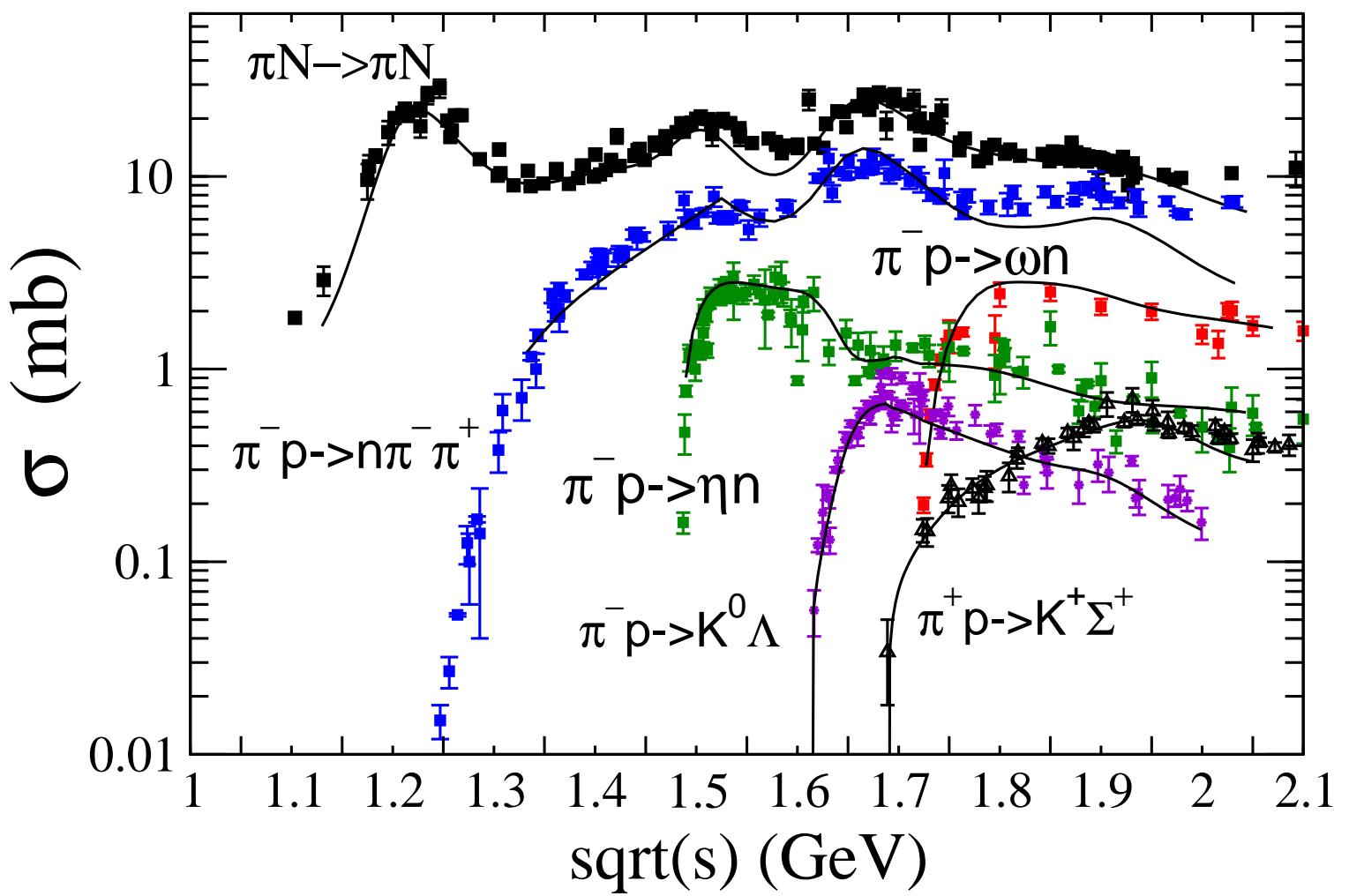


# Models for Extracting $N^*$ Parameters from Meson-Baryon Reactions

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**Data of  $\gamma p$  reaction cross sections  
(from Ostrick and Schmieden )**



Data of  $\pi N$  reaction cross sections  
(from V. Shklyar )

Challenge :

Extensive data of electromagnetic production of  
 $\pi$ ,  $\eta$ ,  $K$ ,  $\omega$ ,  $\phi$ , and  $\pi\pi N$  ( $\rho N$ ,  $\pi\Delta$ )



Extract properties of nucleon resonances ( $N^*$ )



Understand non-perturbative QCD :

- Confinement of constituent quarks
- Chiral dynamics of meson cloud of baryons

Tasks :

- Perform Amplitude Analyses of data

→

Extract  $N^*$  parameters

- Develop Dynamical Reaction Models

→

Interpret  $N^*$  parameters in terms of QCD :

- Hadron Models (now )
- Lattice QCD (near future )

Status :

- In the  $\Delta$  region
  - Well developed
  - Amplitude Analyses and Dynamical Reaction Models are complementary

Example:  $\gamma N \rightarrow \Delta$  M1 transition

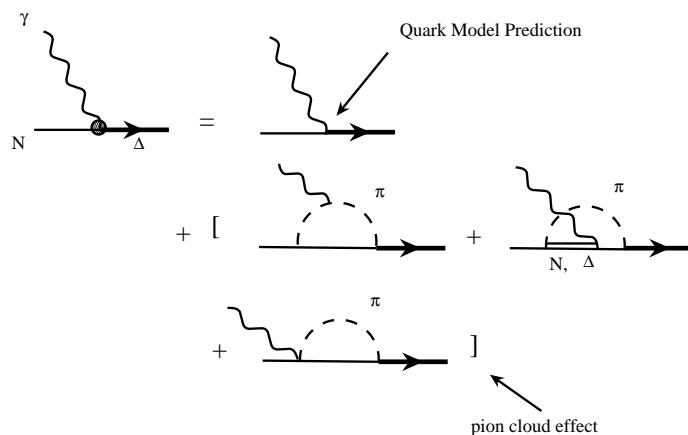
- All amplitude analyses :  $G_M(0) = 3.50 \pm 0.2$
- Disagree with the quark model :

$$\frac{G_M^{\text{Exp.}}(0)}{G_M^{\text{Q.M.}}(0)} \sim 1.4$$

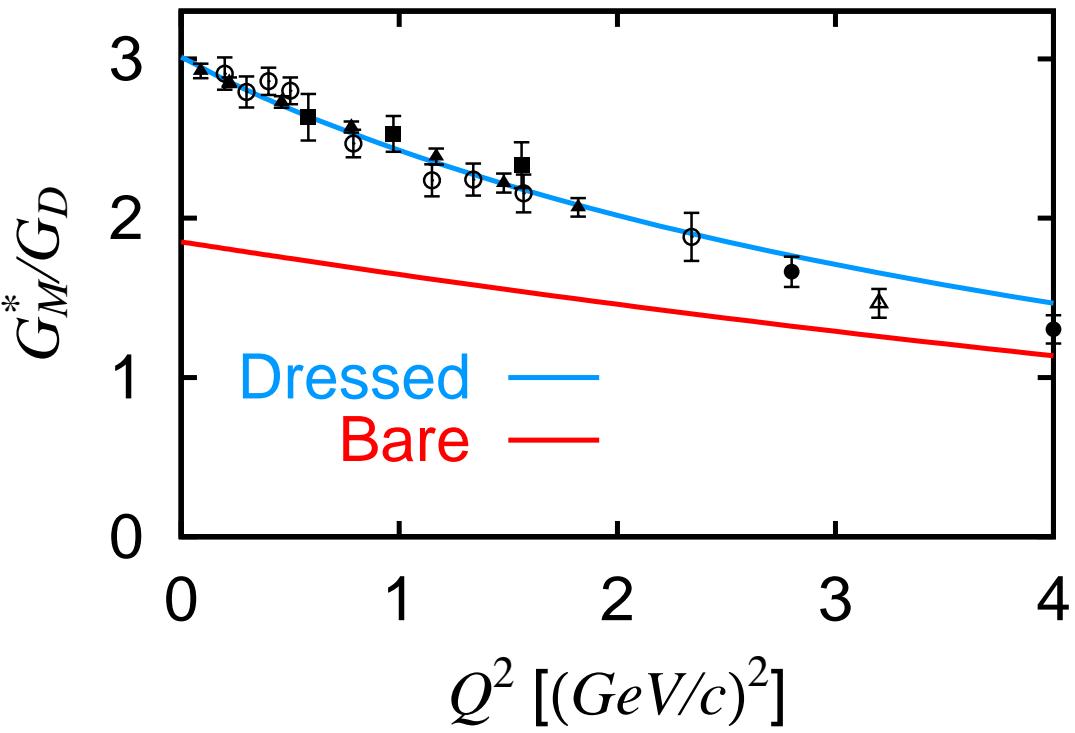
Solution: Develop Dynamical Models

→

Find: due to meson cloud



$$\gamma N \rightarrow \Delta \text{ Magnetic Dipole} \quad G_M(Q^2)$$

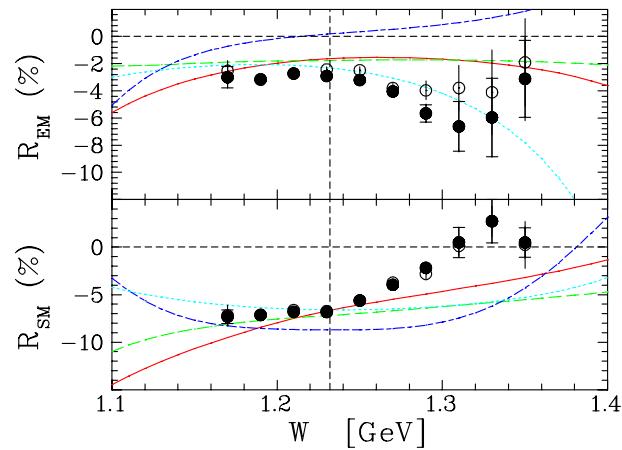
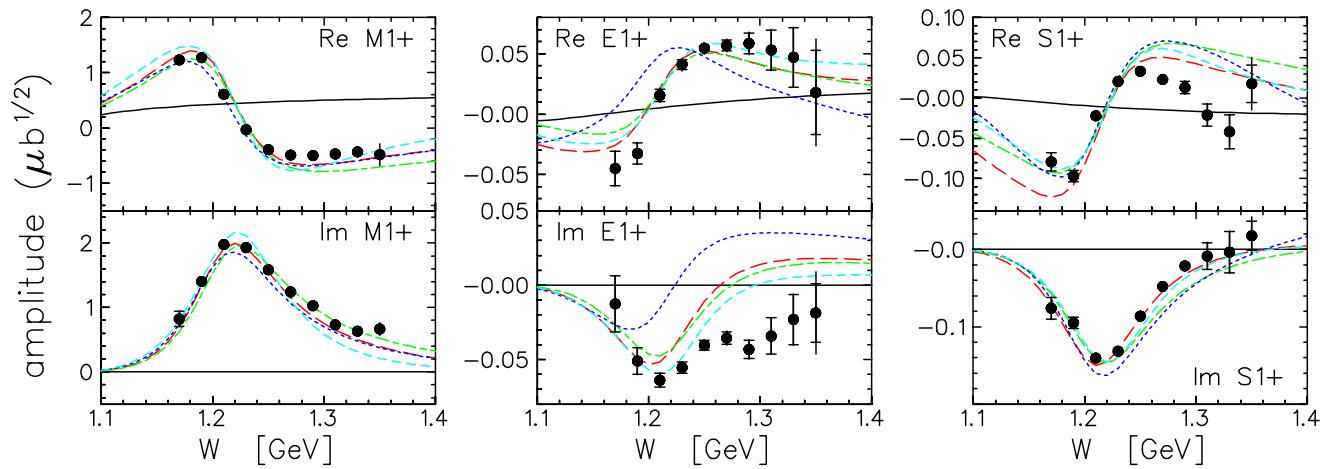


- Pion cloud has a very large effect on  $G_M$

Recent progress in the  $\Delta$  region :

- 16 response functions of  $p(\vec{e}, e\vec{p})$  at  $Q^2 = 1$  (GeV/c)<sup>2</sup> have been obtained at JLab  
(allow almost model independent amplitude analysis)
- LQCD calculations of  $N$ - $\Delta$  form factors are available
- Low  $Q^2$  data have been obtained at JLab, Mainz, MIT-Bates  
( reveal  $Q^2$ -evolution of meson cloud effects on  $N$ - $\Delta$  )
- Data of  $d(\vec{\gamma}, \pi N)N$  have been obtained at LEGS  
(will provide  $\gamma n \rightarrow \pi N$  multipoles :A. Sandorfi's talk )

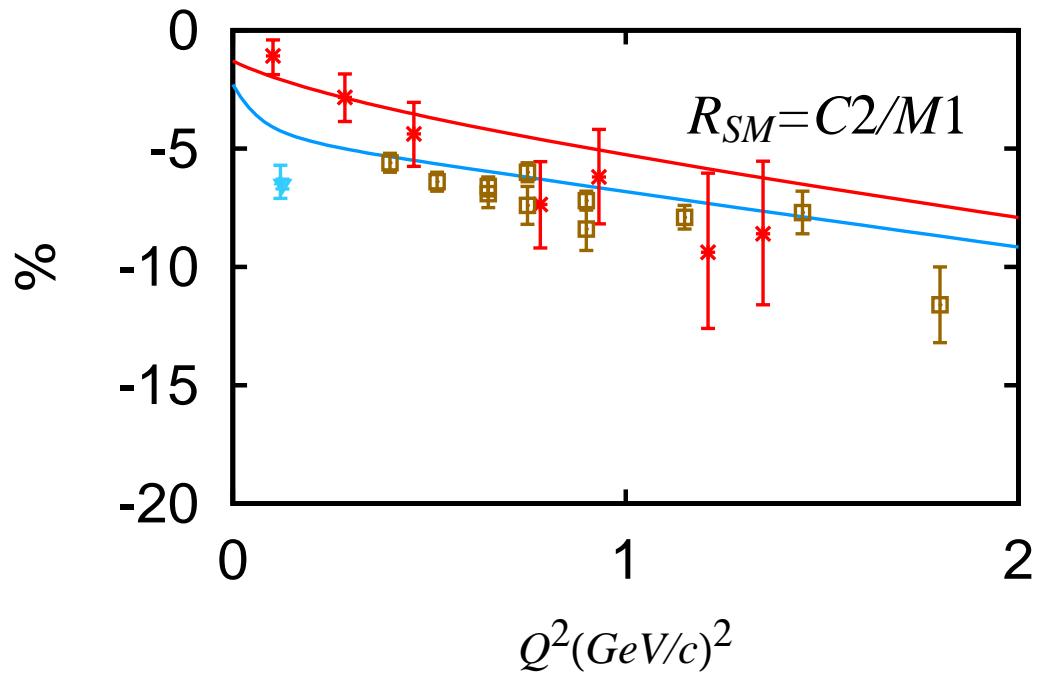
## Model Independent Amplitudes at $Q^2 = 1$ ( $\text{GeV}/c$ )<sup>2</sup> (J. Kelly et al. (2005) )



Curves : SL, DMT, MAID, SAID

## Recent Results from LQCD

$R_{SM}$  : measure deformation of  $N$  or  $\Delta$



red-bar : LQCD of C. Alexandrou et al. (2005)

red curve : Bare f.f of SL Model

blue curve : Dressed f.f. of SL Model

- In the second and third resonance regions:

Open channels:  $\eta N$ ,  $\pi\pi N$  ( $\pi\Delta$ ,  $\rho N$ ),  $\omega N$ ,  $KN \dots$

→

Need to develop coupled-channel approaches

This talk : Review current developments

## Outline

- Introduce a reaction formulation :
  - derive and compare  
current models of meson production reactions
- Review the status of the coupled-channel analyses
- Concluding Remarks

## Reaction Theories

- Based on Hamiltonian or Bethe-Salpeter Equations :

$$T(E) = V + V \frac{1}{E - H_0 + i\epsilon} T(E)$$

$V$  = *interactions*

Can be used to derive

- Unitary Isobar Models :  
MAID  
Jlab/Yerevan UIM
- Multi-channel K-matrix models :  
SAID  
Giessen model, KVI model  
Kent State University ( KSU )
- Carnegie-Mellon Berkeley (CMB) Model
- Dynamical Reaction models  
Juelich, SL, DMT, Ohio-Utrecht . . .  
Chiral SU(3) models

- Based on Dispersion Relations :

$$Re A^I(s, t) = B^I + \frac{1}{\pi} P \int_{s_{thr}}^{\infty} \left[ \frac{1}{s' - s} + \frac{\epsilon^I \xi_i}{s' - u} \right] Im A^I(s', t)$$

Recent works :

- Constraint on  $\pi N$  amplitudes in SAID  
R.Arndt, I. Strakovsky, R. Workman  
and collaborators (1996, 2004)
- $\gamma N \rightarrow \pi N$   
O. Hanstein, D. Drechsel, and L. Tiator (1998)
- $\gamma N \rightarrow \pi N, \eta N$   
I. Aznauryan (1998, 2003)

Will not be covered in this talk

## Derivations of Models

- Define  $K$  operator:

$$T(E) = V + V \frac{1}{E - H_0 + i\epsilon} T(E)$$

→

$$T(E) = V + V \left[ \frac{P}{E - H_0} - i\pi\delta(E - H_0) \right] T(E)$$

$$K(E) = V + V \frac{P}{E - H_0} K(E)$$

$P$  : the principal-value integration.

→

$$T(E) = K(E) - T(E)[i\pi\delta(E - H_0)]K(E)$$

→

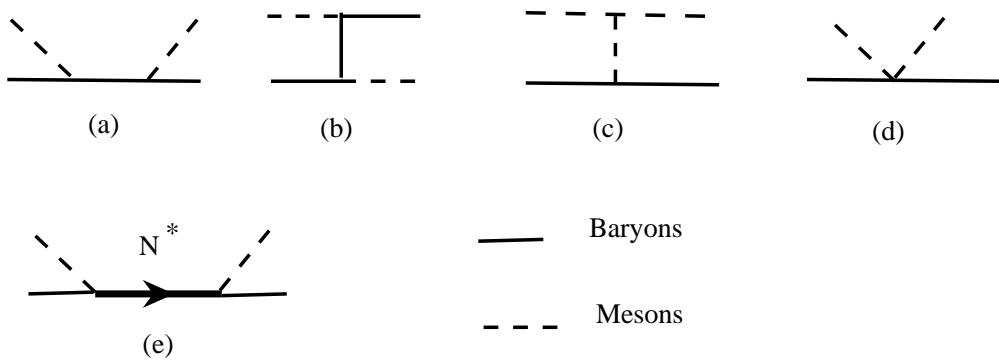
Lead to **on-shell** relations between  $T$  and  $K$

- Define interactions :

$$V = v^{bg} + v^R$$

Non-resonant term:  $v^{bg}$

**Resonant term:**  $v^R = \frac{\Gamma_i^\dagger \Gamma_i}{E - M_{N_i^*}}$



## Approaches :

- Start with  $V = v^{bg} + v^R$  :

$$T_{a,b}(k_a, k_b, E) = V_{a,b}(k_a, k_b) + \sum_c \int dk \frac{V_{a,c}(k_a, \mathbf{k}) T_{c,b}(\mathbf{k}, k_b)}{E - E_{Mc}(\mathbf{k}) - E_{Bc}(\mathbf{k}) + i\epsilon}$$

$$a, b = \pi N, \gamma N, \eta N, \omega N, KY, \rho N, \pi\Delta (\pi\pi N)$$

- Need off-shell information
- Equations for Dynamical Models

- Start with  $K$  matrix:

A matrix relation:

$$T_{a,b}(E) = \sum_c [(1 + iK(E))^{-1}]_{a,c} K_{c,b}(E)$$

$$a, b = \pi N, \gamma N, \eta N, \omega N, KY, \rho N, \pi\Delta (\pi\pi N)$$

- Need only on-shell information
- Equations for K-matrix Models

## Derivations

- Unitary Isobar Model (**UIM**) :

- start with **K** matrix
- channels :  $\gamma N$ ,  $\pi N$  (**or**  $\eta N$ )

→

$\gamma N \rightarrow \pi N$  amplitude :

$$\begin{aligned} T_{\pi N, \gamma N} &= \frac{1}{1 + iK_{\pi N, \pi N}(E)} K_{\pi N, \gamma N}(E) \\ &= e^{i\delta_{\pi N}} \cos \delta_{\pi N} K_{\pi N, \gamma N}(E) \\ &\sim e^{i\delta_{\pi N}} \cos \delta_{\pi N} V_{\pi N, \gamma N} \end{aligned}$$

$V_{\pi N, \gamma N}$  = Tree-diagrams

$\delta_{\pi N}$  :  $\pi N$  phase shifts

→

Satisfy Watson Theorem in  $W < 1.3$  GeV

– Mainz and Jlab/Yerevan UIM :

1. Include of  $N^*$  by using Walker's parameterization
2. Unitarize the total amplitude

→

$$T_{\pi N, \gamma N}(\text{UIM}) = e^\delta \cos \delta [v_{\pi N, \gamma N}^{bg}] + \sum_{N_i^*} T_{\pi N, \gamma N}^{N_i^*}(W)$$
$$T_{\pi N, \gamma N}^{N_i^*}(E) = f_{\pi N}(W) \frac{\Gamma^{tot} M_i e^{i \Phi_i}}{M_i^2 - W^2 - i M_i \Gamma^{tot}} A_{\gamma N}(W)$$

$\Phi_i$  : Unitarization Phase

**Results** from MAID and JLab/Yerevan UIM :

1. Successful in extracting  $\Delta$  parameters
2. Can fit pion production data up to  $W = 2$  GeV.

**Comments :**

Coupled-channel effects are not treated explicitly

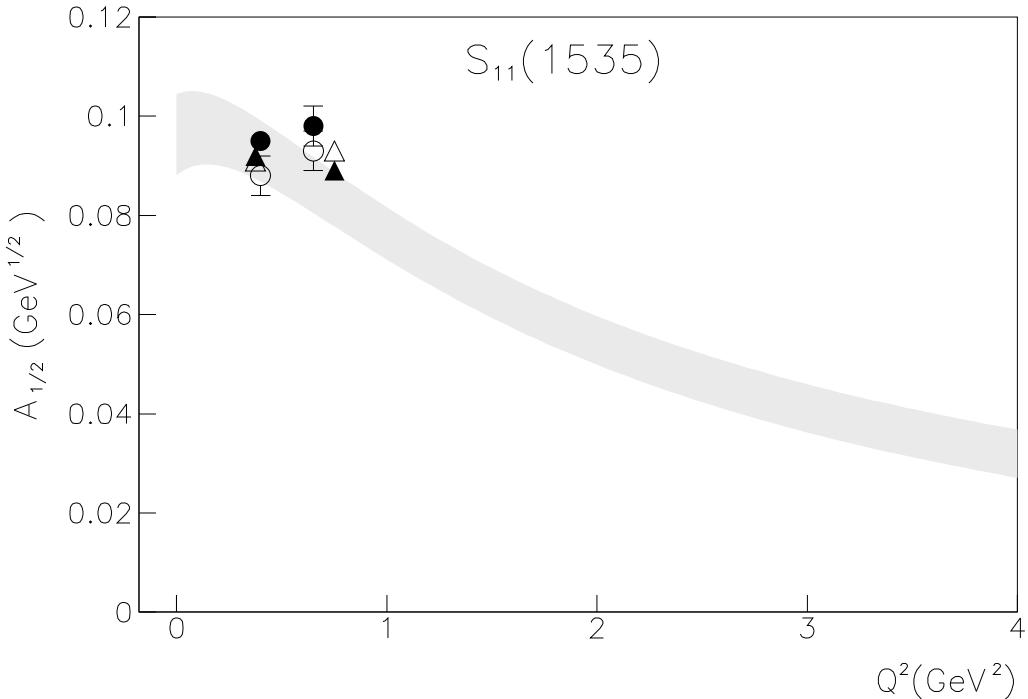
[  $\gamma N \rightarrow (\pi\Delta, \rho N \dots) \rightarrow \pi N$  is neglecte ]

→

The extracted  $N^*$  parameters need to be verified

Jlab/Yerevan UIM **global fits** to  $p\pi^0$ ,  $n\pi^+$ , and  $p\eta$

→ Obtain **highly constrained**  $\gamma N \rightarrow N^*$  form factors



More will be given by I. Aznauryan's talk

Recent results from **MAID** :

$\gamma N \rightarrow N^*$  form factors have been extracted  
(L. Tiator's talk )

- Multi-channel K-matrix models

– SAID :

Consider  $\gamma N, \pi N, \pi\Delta$  (all inelastic channels )

$\rightarrow$

$$T_{\gamma N, \pi N}(\text{SAID}) = A_I(1 + iT_{\pi N, \pi N}) + A_R T_{\pi N, \pi N}$$

$$\begin{aligned} A_I &= K_{\gamma N, \pi N} - \frac{K_{\gamma N, \pi\Delta} K_{\pi N, \pi N}}{K_{\pi N, \pi\Delta}} \\ A_R &= \frac{K_{\gamma N, \pi\Delta}}{K_{\pi N, \pi\Delta}} \end{aligned}$$

Actual analysis:

$$A_I = v_{\gamma N, \pi N}^{bg} + \sum_{n=0}^M \bar{p}_n z Q_{l_\alpha+n}(z)$$

$$A_R = \frac{m_\pi}{k_0} \left( \frac{q_0}{k_0} \right)^{l_\alpha} \sum_{n=0}^N p_n \left( \frac{E_\pi}{m_\pi} \right)^n$$

$\bar{p}_n, p_n$ : fitting parameters

$N^*$  parameters are extracted by fitting the resulting amplitudes to a Breit-Wigner parameterization at  $W \rightarrow M^*$

**Results** from SAID :

1. Determine  $\pi N \rightarrow \pi N$ ,  $\gamma N \rightarrow \pi N$  amplitudes
2. extract  $N^*$  parameters

**Comments** :

Coupled-channel effects are not treated explicitly

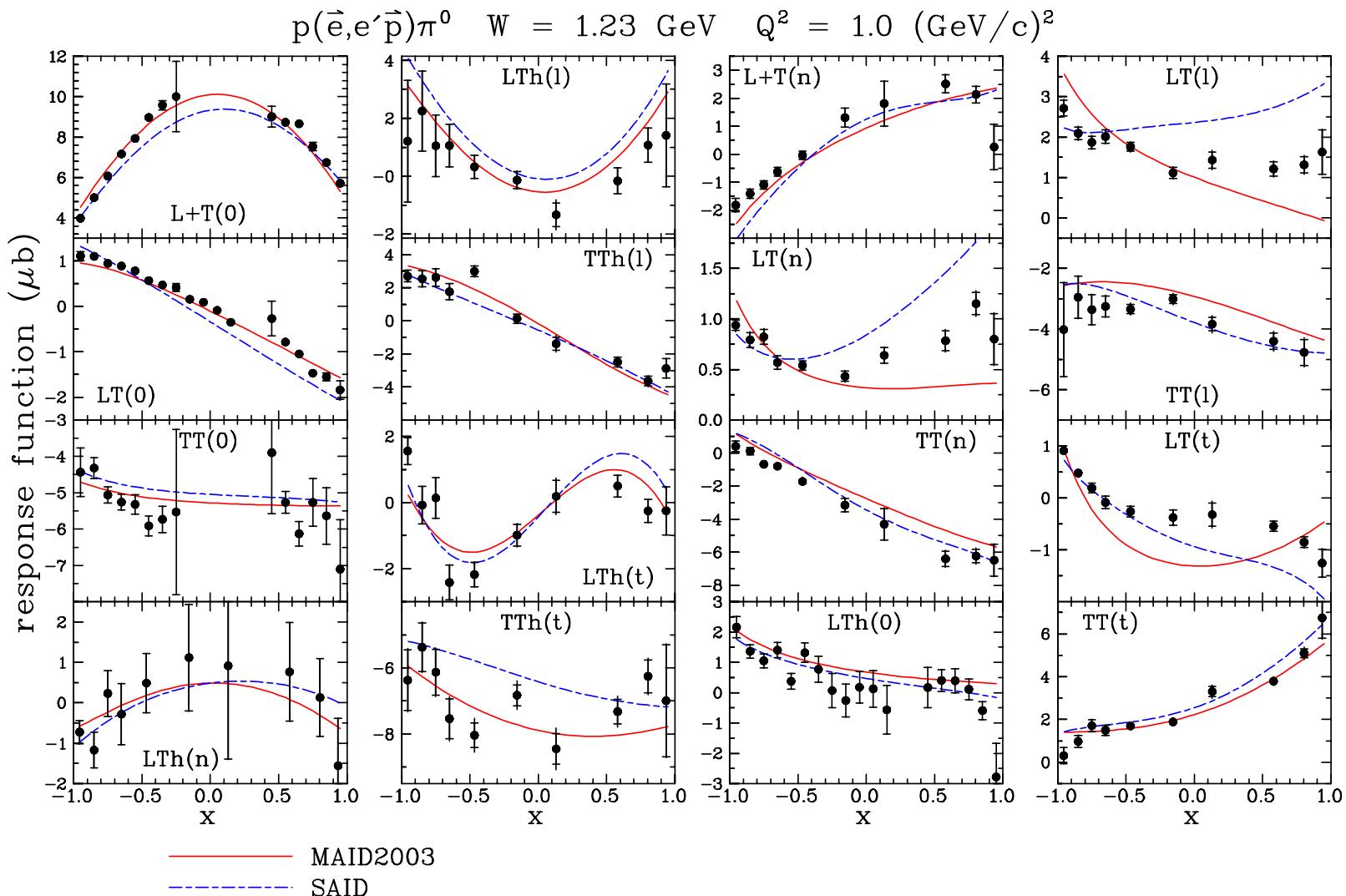
[  $\gamma N \rightarrow (\pi\Delta, \rho N \dots) \rightarrow \pi N$  is neglected ]

→

The extracted  $N^*$  parameters need to be **verified**

## SAID (2005)

$$T_{\gamma N, \pi N}(\text{SAID}) = A_I(1 + iT_{\pi N, \pi N}) + A_R T_{\pi N, \pi N} \\ + (C + iD)(Im T_{\pi N, \pi N} - T_{\pi N, \pi N}^2)$$



**Data :**

16 response functions (J. Kelly et al. (2005))

– Giessen Model and KVI Model :

Set  $K \rightarrow V =$  Tree-diagrams

→

$$T_{a,b}(E) = \sum_c [(1 + iV(E))^{-1}]_{a,c} V_{c,b}(E)$$

Comments :

Higher-order and off-shell effects are neglected

(Note :  $K = V + V \frac{P}{E - H_0} K$  )

→

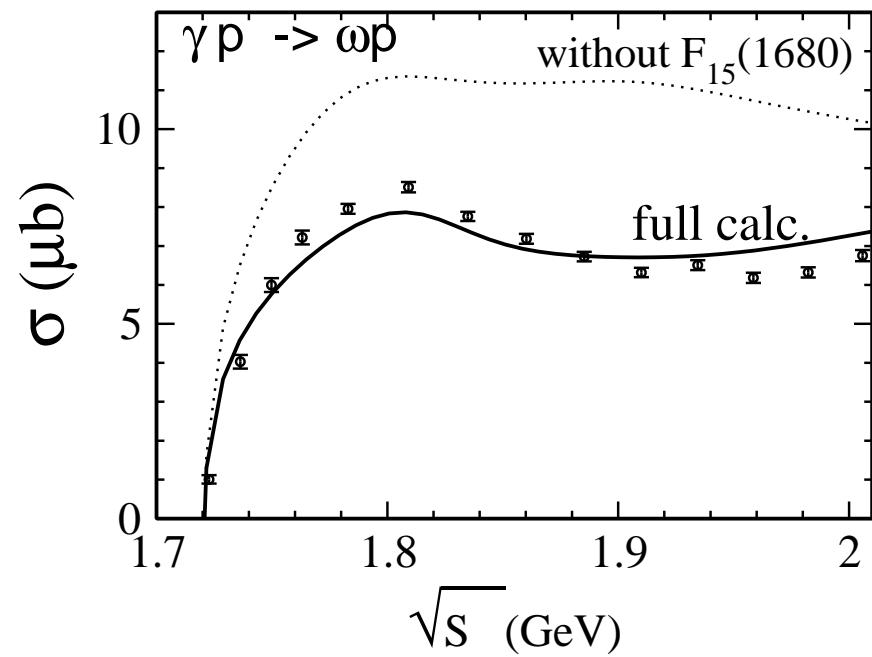
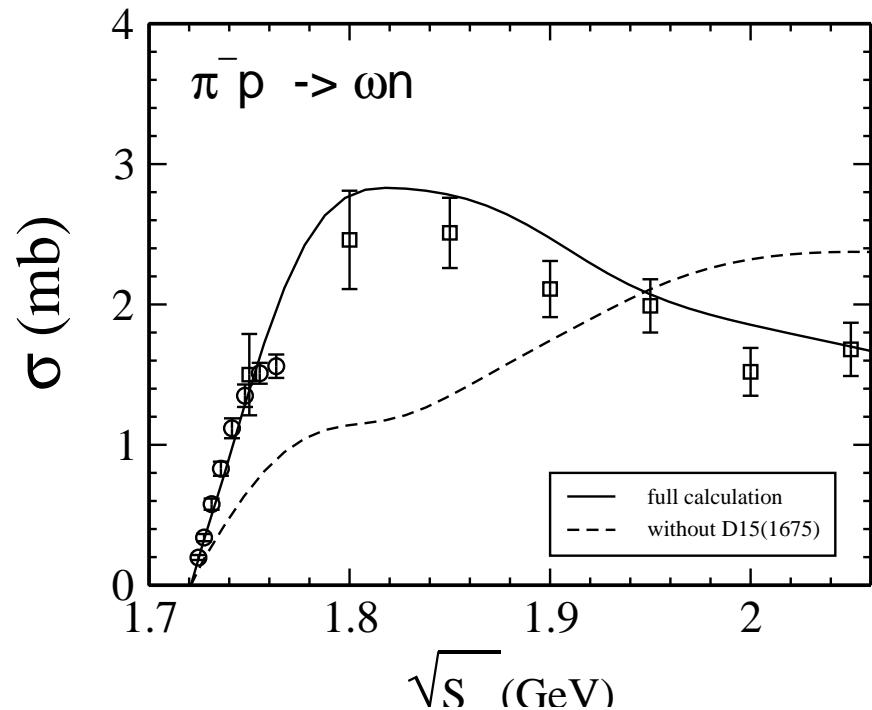
The extracted  $N^*$  parameters can not be interpreted in terms of quark models and/or LQCD

(Lesson from the study of  $\Delta$  )

Recent results of Giessen Model  
(V. Shklyar et al. (2005) ):

1. Include  $\gamma N$ ,  $\pi N$ ,  $2\pi N$ ,  $\eta N$ , and  $\omega N$ .
2. Find **evidence** for  $D_{15}(1675)$  in  $\pi N \rightarrow \omega N$
3. Find **evidence** for  $(F_{15}(1680))$  in  $\gamma N \rightarrow \omega N$

## Recent results from Giessen Model (2005)



## Recent Results of KVI Model (A. Usov and O. Scholten (2005) )

- Include  $\pi N$ ,  $\gamma N$ ,  $K\Lambda$ ,  $K\Sigma$ ,  $\phi N$ ,  $\eta N$
- Fit the data of  $\gamma p \rightarrow K^+\Lambda$ ,  $K^+\Sigma^0$ ,  $K^0\Sigma^+$   
(Show large coupled-channel effects)

Comments :

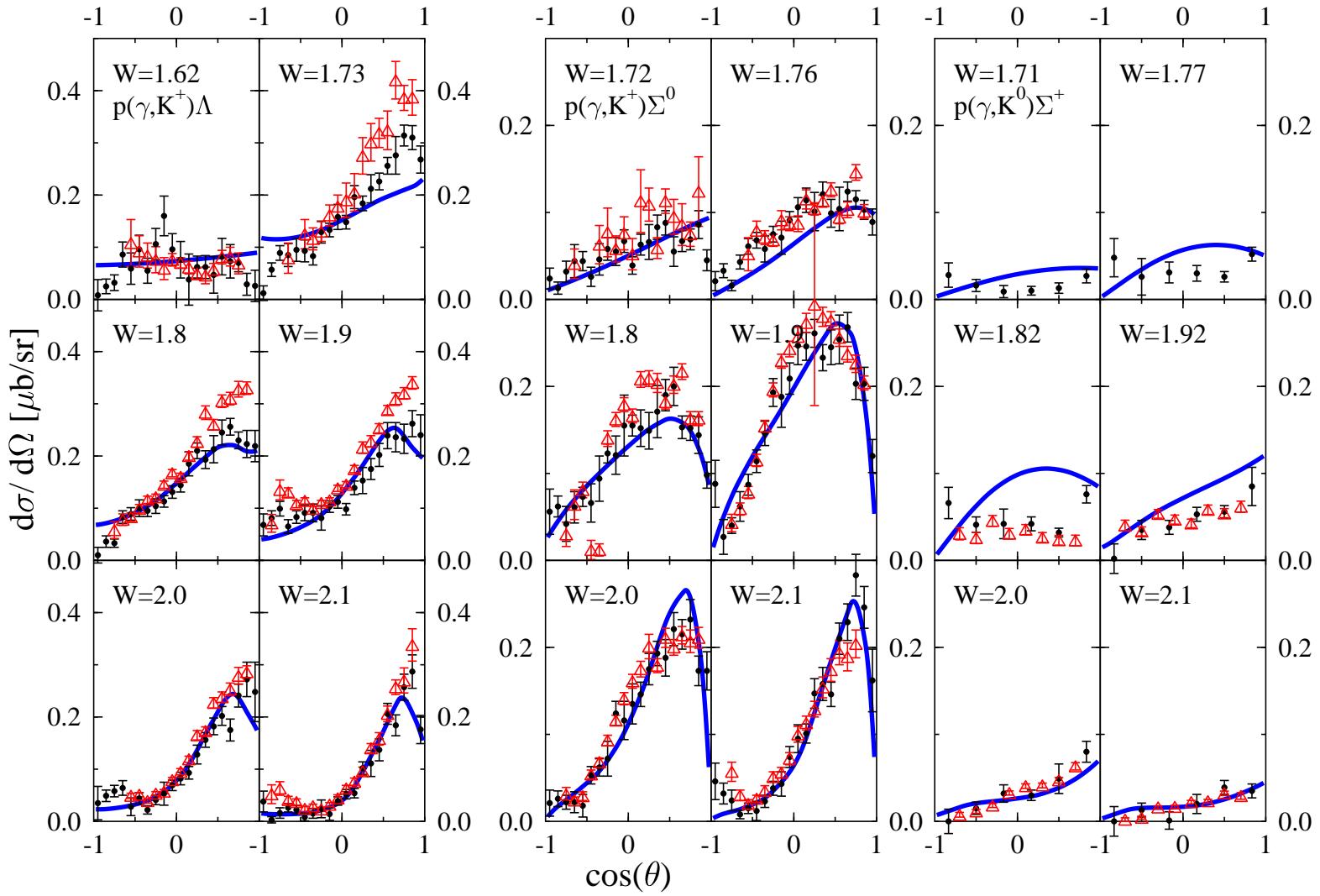
$2\pi N$  channel is not included

(Note :  $\sigma_{\gamma N \rightarrow \pi\pi N} \gg \sigma_{\gamma N \rightarrow KY}$  )

→

The extracted  $N^*$  parameters need to be verified

## $\gamma p \rightarrow KY$ results of KVI Model (2005)



More will be given in O. Scholten's talk

For deriving:

- Carnegie-Mellon Berkeley (**CMB**) Model
- Kent State University (**KSU**) model
- **Dynamical models**

Apply **two-potential scattering formulation**

for  $V = v^{bg} + \frac{\Gamma_{N^*}^\dagger \Gamma_{N^*}}{E - M_{N^*}^0}$   
→

$$T(E) = t^{bg}(E) + \frac{\bar{\Gamma}_{N^*}^\dagger(E) \bar{\Gamma}_{N^*}(E)}{E - M_{N^*}^0 - \Sigma_{N^*}(E)}$$

$$t^{bg} = v^{bg} + v^{bg} G(E) t^{bg}(E)$$

**Resonance parameters :**

$$\bar{\Gamma}_{N^*} = \Gamma_{N^*} + \Gamma_{N^*} G(E) t^{bg}(E)$$

$$\Sigma_{N^*}(E) = \Gamma_{N^*}^\dagger G(E) \bar{\Gamma}_{N^*}$$

**Main features :**

- Isolate resonant term  $\sim$  **Briet-Wigner form**
- Non-resonant effects on resonance parameters are identified

For **multi** -channel **multi** -resonant case:

$$T_{a,b}(E) = t_{a,b}^{bg}(E) + \sum_{N_i^*, N_j^*} \bar{\Gamma}_{N_i^*, a}^\dagger(E) [\hat{G}(E)]_{i,j} \bar{\Gamma}_{N_j^*, b}(E)$$

$$t_{a,b}^{bg} = v_{a,b}^{bg} + \sum_c v_{a,c}^{bg} G_c(E) t_{c,b}^{bg}(E)$$

$$\bar{\Gamma}_{N^*, a} = \Gamma_{N^*, a} + \sum_b \Gamma_{N^*, b} G_b(E) t_{b,a}^{bg}$$

$$[\hat{G}(E)^{-1}]_{i,j}(E) = (E - M_{N_i^*}^0) \delta_{i,j} - \Sigma_{i,j}(E)$$

$$\Sigma_{i,j}(E) = \sum_a \Gamma_{N^*, a}^\dagger G_a(E) \bar{\Gamma}_{N_j^*, a}$$

$$\mathbf{a}, \mathbf{b} = \pi N, \gamma N, \eta N, \omega N, KY, \sigma N, \rho N, \pi\Delta (\pi\pi N)$$

- Carnegie-Mellon Berkeley (CMB) Model

**Set :**  $v_{a,b}^{bg}(E) = \frac{\Gamma_{\text{L},a}^\dagger \Gamma_{\text{L},b}}{E - M_{\text{L}}} + \frac{\Gamma_{\text{H},a}^\dagger \Gamma_{\text{H},b}}{E - M_{\text{H}}}$

→

$$V = v^{bg} + v^R = \sum_{i=N_i^*, \text{L}, \text{H}} \frac{\Gamma_{i,a}^\dagger \Gamma_{i,b}}{E - M_i} = \text{Separable}$$

→

$$T_{a,b}(E) = \sum_{i,j} \Gamma_{i,a}^\dagger G_{i,j}(E) \Gamma_{j,b}$$

$$G(E)_{i,j}^{-1} = (E - M_i^0) \delta_{i,j} - \Sigma_{i,j}(E)$$

$$\Sigma_{i,j}(E) = \sum_a \int k^2 dk \frac{\Gamma_{i,a}^\dagger(k) \Gamma_{j,a}(k)}{E - E_{M_a}(k) - E_{B_a}(k) + i\epsilon}$$

With appropriate variable changes :  $s = E^2$

→

CMB's dispersion relations :

$$\Sigma_{i,j}(s) = \sum_c \gamma_{i,c} \Phi_c(s) \gamma_{j,c}$$

$$Re[\Phi_c(s)] = Re[\Phi_c(s_0)] + \frac{s - s_{th,c}}{\pi} \int_{s_{th}}^{\infty} \frac{Im[\Phi_c(s')]}{(s' - s)(s' - s_0)} ds'$$

→

CMB model is analytic

**Recent** applications/extensions of CMB model :

- **Zagreb** : M. Batinic, A. Svarc and collaborators

Consider **three** channels :  $\pi N$ ,  $\eta N$ ,  $\sigma(\pi\pi)N$

- **PITT-ANL** : T. Varana, S. Dytman, T.-S. H. Lee

Consider up to **eight** channels:

$\pi N$ ,  $\eta N$ ,  $\pi\Delta$ ,  $\rho N$ ,  $\sigma(\pi\pi)N$ ,  $\pi N^*(1440)$ ,  $K\Lambda$ ,  $\gamma N$

- **FSU-PITT** :

A. Kiswandhi, S. Capstick, and S. Dytman

Investigate **model-dependence** in  $S_{11}$  channel

Results:

- $N^*$  in  $S_{11}$  channel is better understood
- The interplay between channel coupling and  $N^*$  excitation has been better understood
- Some extracted  $N^*$  parameters are significantly different from PDG values

Current effort (A. Kiswandhi, S. Capstick ) :

Approach is being developed to replace

$$v_{a,b}^{bg}(E) = \frac{\Gamma_{\textcolor{red}{L},a}^\dagger \Gamma_{\textcolor{red}{L},b}}{E - M_{\textcolor{red}{L}}} + \frac{\Gamma_{\textcolor{red}{H},a}^\dagger \Gamma_{\textcolor{red}{H},b}}{E - M_{\textcolor{red}{H}}}$$

by dynamical models

- Kent State University (**KSU**) model

Start with

$$T(E) = t^{bg}(E) + \frac{\bar{\Gamma}_{N^*}^\dagger \bar{\Gamma}_{N^*}}{E - M_{N^*}^0 - \Sigma_{N^*}(E)}$$

One can derive exactly the **distorted-wave** form

$$\begin{aligned} S(E) &= 1 + 2iT(E) \\ &= \omega^{(+)} R(E) \omega^{(+)} \end{aligned}$$

where

$$\begin{aligned} \omega^{(+)} &= 1 + G(E)t^{bg}(E) \\ R(E) &= 1 + 2iT^R(E) \\ T^R(E) &= \frac{\Gamma_{N^*}^\dagger(E) \Gamma_{N^*}(E)}{E - M_{N^*}^0 - \Sigma_{N^*}(E)} \end{aligned}$$

**KSU** separable parameterization:

$$T^R(E) = \frac{K}{1 + iK}$$

$$K_{ij} = \sum_\alpha \tan \delta_\alpha f_{i\alpha} f_{j\alpha}$$

$$\omega^{(+)} = B_1 B_2 \cdots B_n$$

$$B_i \sim e^{iX\Delta_i}$$

## Recent Results from KSU Model

- Fit  $S = -1$  amplitudes with Channels :

$\bar{K}n, \pi\Lambda, \pi\Sigma, \pi\Sigma^*(1385), \pi\Sigma^*(1520), \bar{K}\Delta, \bar{K}^*N, \eta N$

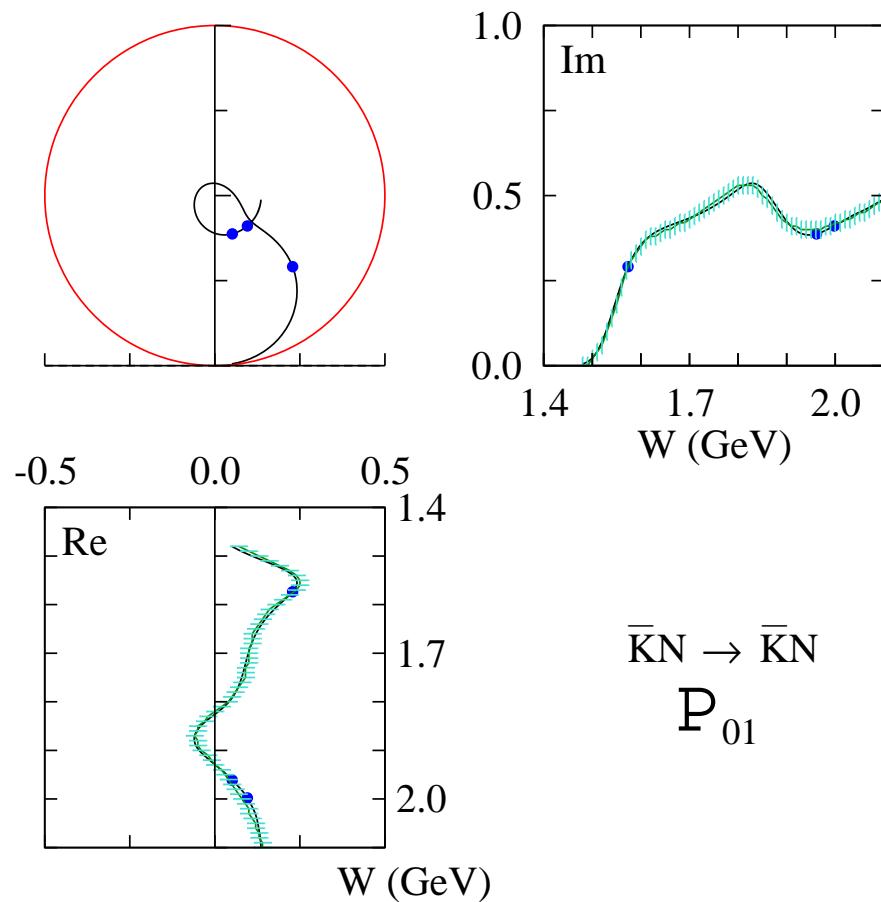
→

Extract  $\Lambda^*$  and  $\Sigma^*$  in  $W = 1560 - 1685$  MeV

- Analyze data of Crystal Ball Collaboration

$K^-p \rightarrow$  neutrals ( $\bar{K}^0n, \pi^0\Lambda, \pi^0\Sigma^0 \dots$ )

## KSU fits to $\bar{K}N \rightarrow \bar{K}N$ (2005)



## Dynamical Models

Two equivalent approaches:

- Solve dynamical equations with  $V = v^{bg} + v^R$  directly :

$$T_{a,b}(E) = V_{a,b} + \sum_c V_{a,c} G_c(E) T_{c,b}(E)$$

$$a, b, c = \pi N, \gamma N, \eta N, \pi \Delta \dots$$

Recent works :

- Juelich Model :  $\pi N$
- Fuda et al. :  $\pi N, \gamma N$
- DMT Model :  $\pi N, \gamma N, \eta N$
- Ohio-Utrecht Model :  $\pi N, \gamma N$
- Chiral SU(3) models :  $KY, \omega N, \gamma N, \pi N$   
(set  $v^R = 0$ )

- Use two-potential formulation to identify resonant mechanism

$$T_{a,b}(E) = t_{a,b}^{bg}(E) + \sum_{N_i^*, N_j^*} \bar{\Gamma}_{N_i^*, a}^\dagger [D^{-1}(E)]_{i,j} \bar{\Gamma}_{N_j^*, b}$$

$$t_{a,b}^{bg}(E) = v_{a,b}^{bg} + \sum_c v_{a,c}^{bg} G_c(E) t_{c,b}^{bg}(E)$$

$$\bar{\Gamma}_{N^*, a} = \Gamma_{N^*, a} + \sum_b \Gamma_{N^*, b} G_b(E) t_{b,a}^{bg}(E)$$

## Recent Works

- Sato-Lee Model :  $\pi N, \gamma N$
- Yoshimoto et al. :  $\pi N, \eta N, \pi\Delta$
- Julia-Diaz et al. :  $\gamma N, KY, \pi N$
- Matsuyama, Lee, Sato :  $\gamma N, \pi N, \eta N, \omega N, \pi\pi N$

## Juelich's Coupled-channel Model

O. Krehl, C. Hanhart, S. Krewald, J. Speth (2000)

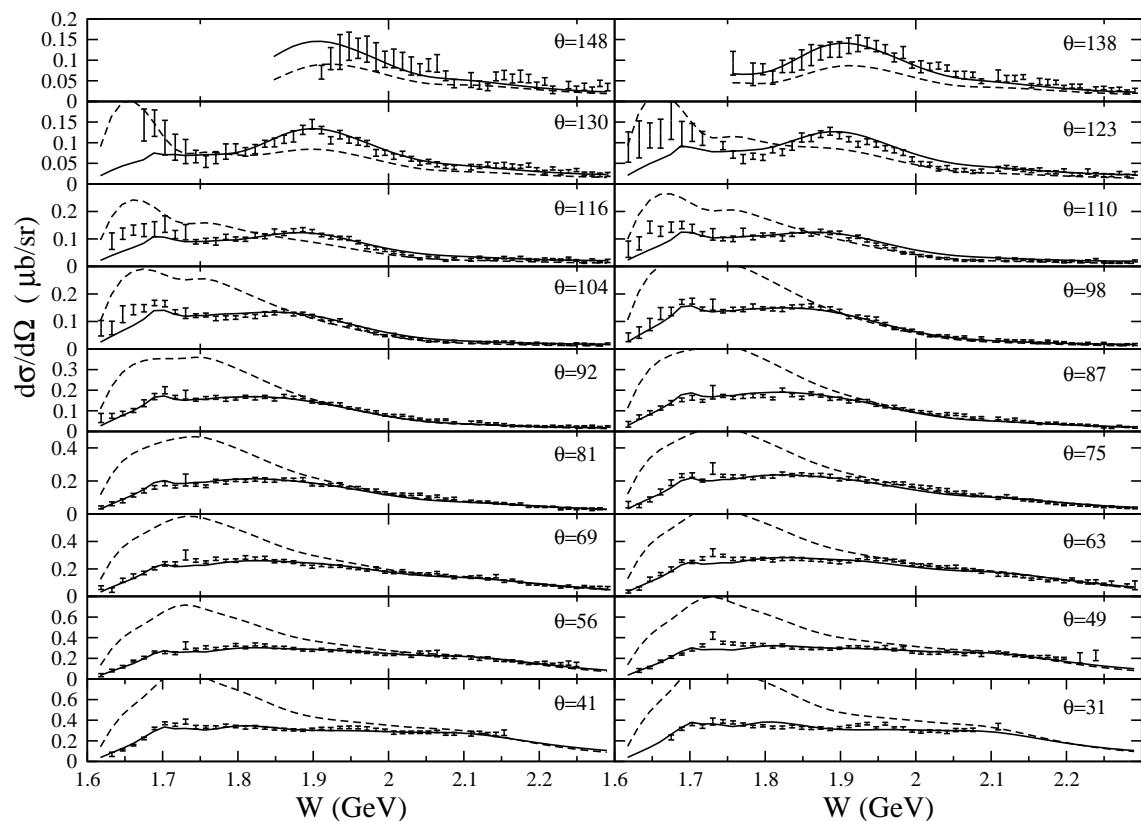
- **Channels** :  $\pi N, \eta N, \sigma N, \pi \Delta, \rho N$ .
- **Main result:**  
 $P_{11}$  is due to meson-baryon coupled-channel effects

Being revised and extended to investigate  $\gamma N \rightarrow \pi N$   
(K. Nakayama et al. (2005) )

## Coupled-channel Model for $KY$ production

B. Julia-Diaz, B.Saghai, F. Tabakin, T.-S. H. Lee (2005)

- **Channels** :  $\gamma N$ ,  $\pi N$ ,  $K\Lambda$ ,  $K\Sigma$
- fit **SAPHIR** and **JLAB** data of  $\gamma N \rightarrow K^+ \Lambda$
- **Main Result** :  
**Large** coupled-channel effects due to  $\pi N$  channel



Recent Results of B. Julia-Diaz et al. (2005)

**More** will be presented in T.-S. H Lee's talk

## DMT Coupled-channel Model

C.-Y Chen, S. Kamalov, S.N. Yang, D. Drechsel, L. Tiator

- **Channels** :  $\gamma N, \pi N, \eta N$
- $\pi\pi N$  effects are assumed to be in  $\Gamma_{N^*} = \Gamma_{1\pi} + \Gamma_{2\pi}$

- **Main Results** :

Need 4  $N^*$  to fit  $S_{11}$  amplitudes

(show **meson cloud** effects on  $\gamma N \rightarrow S_{11}(1535)$ )

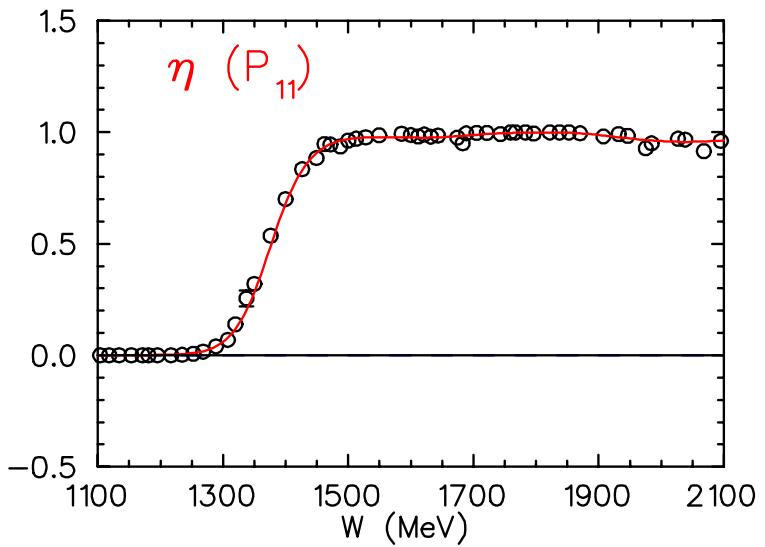
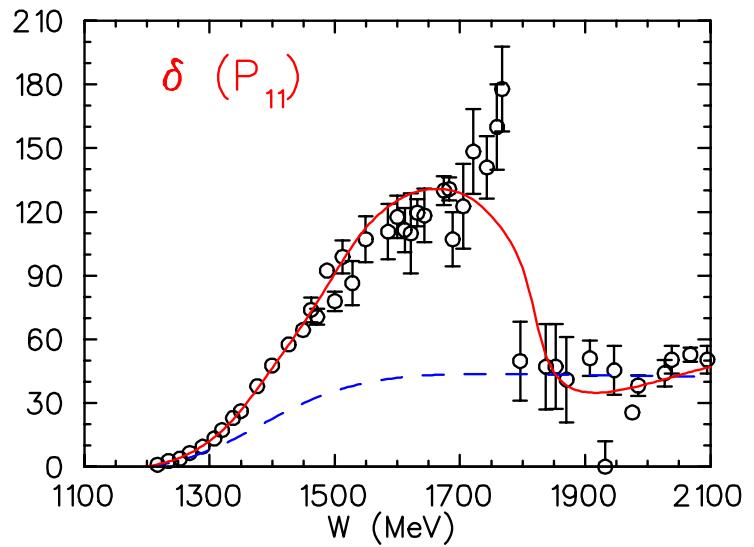
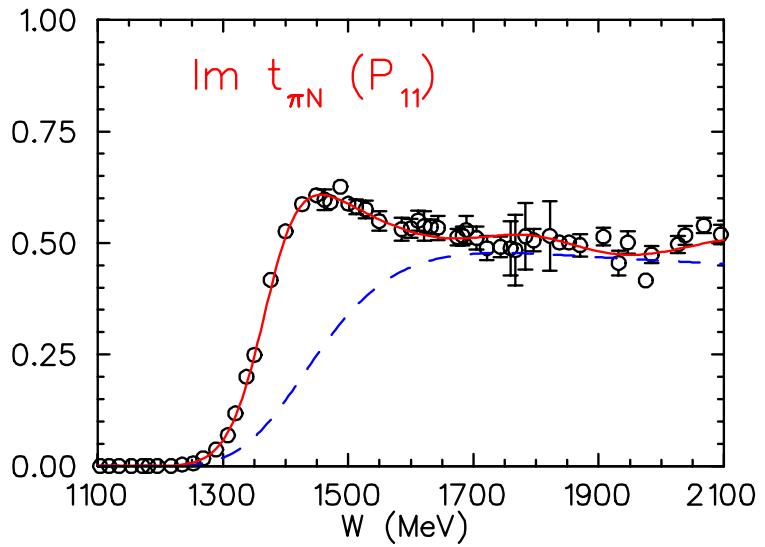
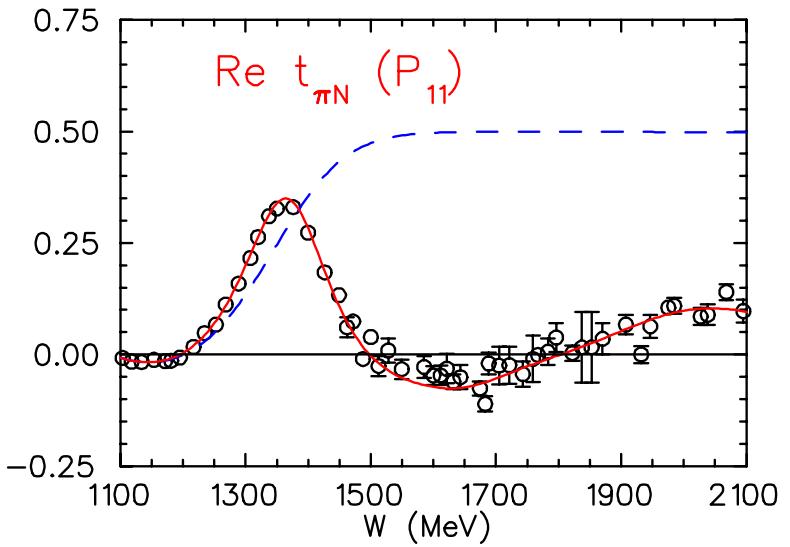
Fit to **all** partial waves will soon be completed

**Comments** :

Need to **justify** its simple treatment of  $\pi\pi N$  channels

(**Note** :  $\sigma_{2\pi N} > \sigma_{\pi N}$  at  $W > 1.5$  GeV)

## Recent Results from DMT Model



## Coupled-channel Model with $\pi\pi N$ channel

A. Matsuyama, T. Sato , T.-S. H. Lee (in progress)

- **Channels** :  $\gamma N, \pi N, \eta N, \omega N, \pi\pi N$  ( $\pi\Delta, \rho N, \sigma N$ )
- Apply second order unitary transformation on  $H$
- Satisfy  $\pi\pi N$  unitarity condition :

$$Im f_{a,a}(\theta = 0) = \sum_b \sigma_{a,b} + \sigma_{a,\pi\pi N}$$

$a, b = \pi N, \gamma N, \eta N, \omega N$  (**stable** particle channels)



Coupled-channel equation with  $\pi\pi N$  cut

Coupled-channel equation with  $\pi\pi N$  cut :

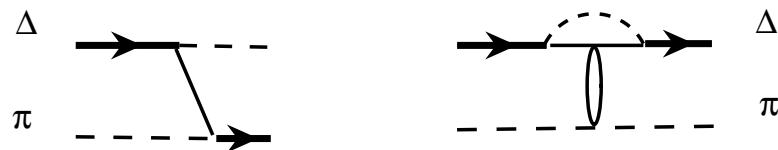
$$X_{a,b}(E) = Z_{a,b}(E) + \sum_c Z_{a,c}(E) G_c(E) X_{cb}(E)$$

$$a, b = \gamma N, \pi N, \eta N, \omega N, \pi\Delta, \rho N, \sigma N$$

$$Z(E) = v^{bg} + Z^{cut}(E)$$

$v^{bg}$  = (tree-diagrams of Chiral Lagrangians )

$Z^{cut}(E)$  :



Apply **Spline** function method :

- solve coupled-channel equations with  $\pi\pi N$  cut

- include  $\pi\pi N$  cut effects **exactly** to calculate



(Note : can not be achieved by **contour rotation** )

- Main Results (2005) :
  - $\gamma N \rightarrow S_{11}(1535)$  :
    1. Meson cloud effect is about 20%
    2. Bare helicity amplitude is close to quark model
  - $\pi\pi N$  unitary cut is crucial in predicting  $\pi\pi N$  production cross sections

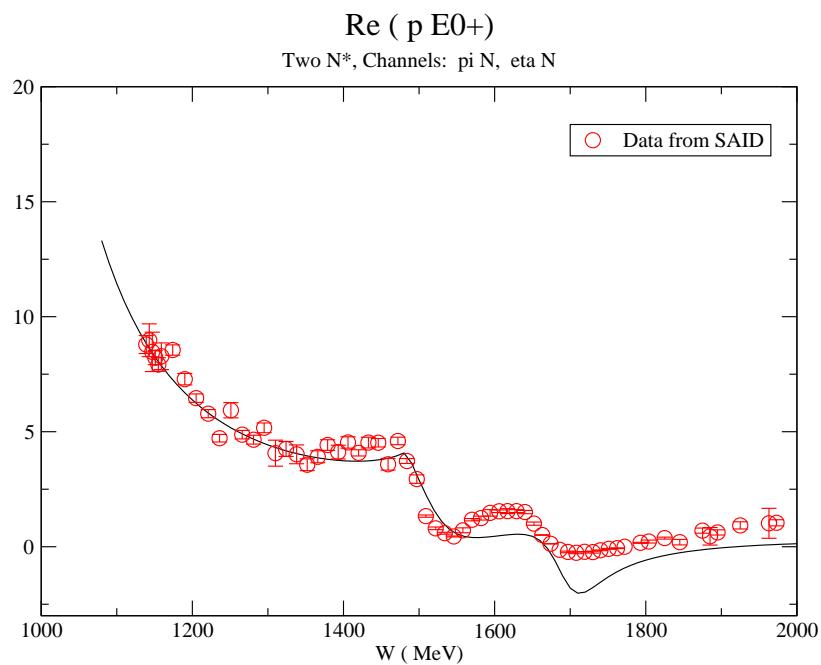
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$\gamma N \rightarrow S_{11}(1535)$		
$\pi\pi N$ model	Dressed	61.24
	Bare	77.64
Capstick		76

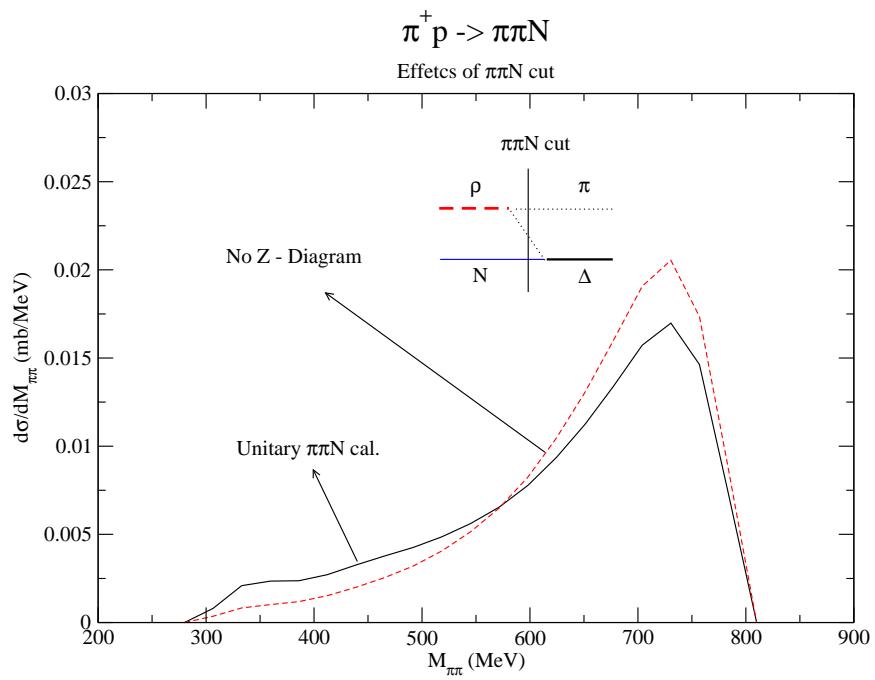
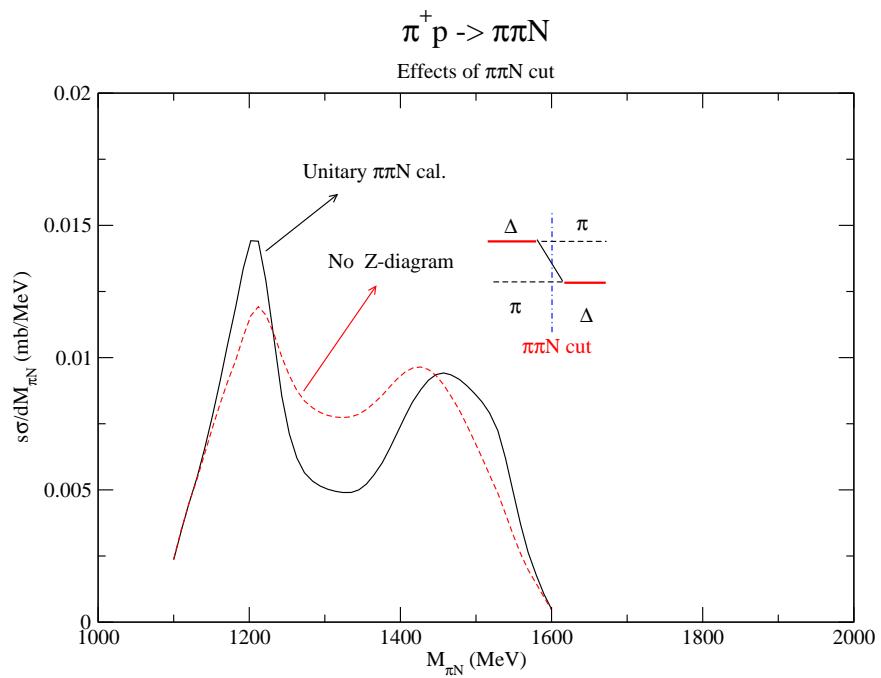
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$$\bar{\Gamma}_{\gamma N, N^*} = \Gamma_{\gamma N, N^*} + \sum_{MB=\pi N, \eta N, \pi\Delta} v_{\gamma N, MB}^{bg} G_{MB} \bar{\Gamma}_{MB, N^*}$$

*Bare*



## Effect of $\pi\pi N$ cut



If coupled-channel and  $\pi\pi N$  cut are neglected

→

- Moscow/JLab Isobar Model of  $\gamma N \rightarrow \pi\pi N$   
(V. Mokeev's talk )
- Amplitude analyses of  $\gamma N \rightarrow \pi\pi N$   
of RPI/JLab

## Chiral $SU(3)$ coupled-channel models

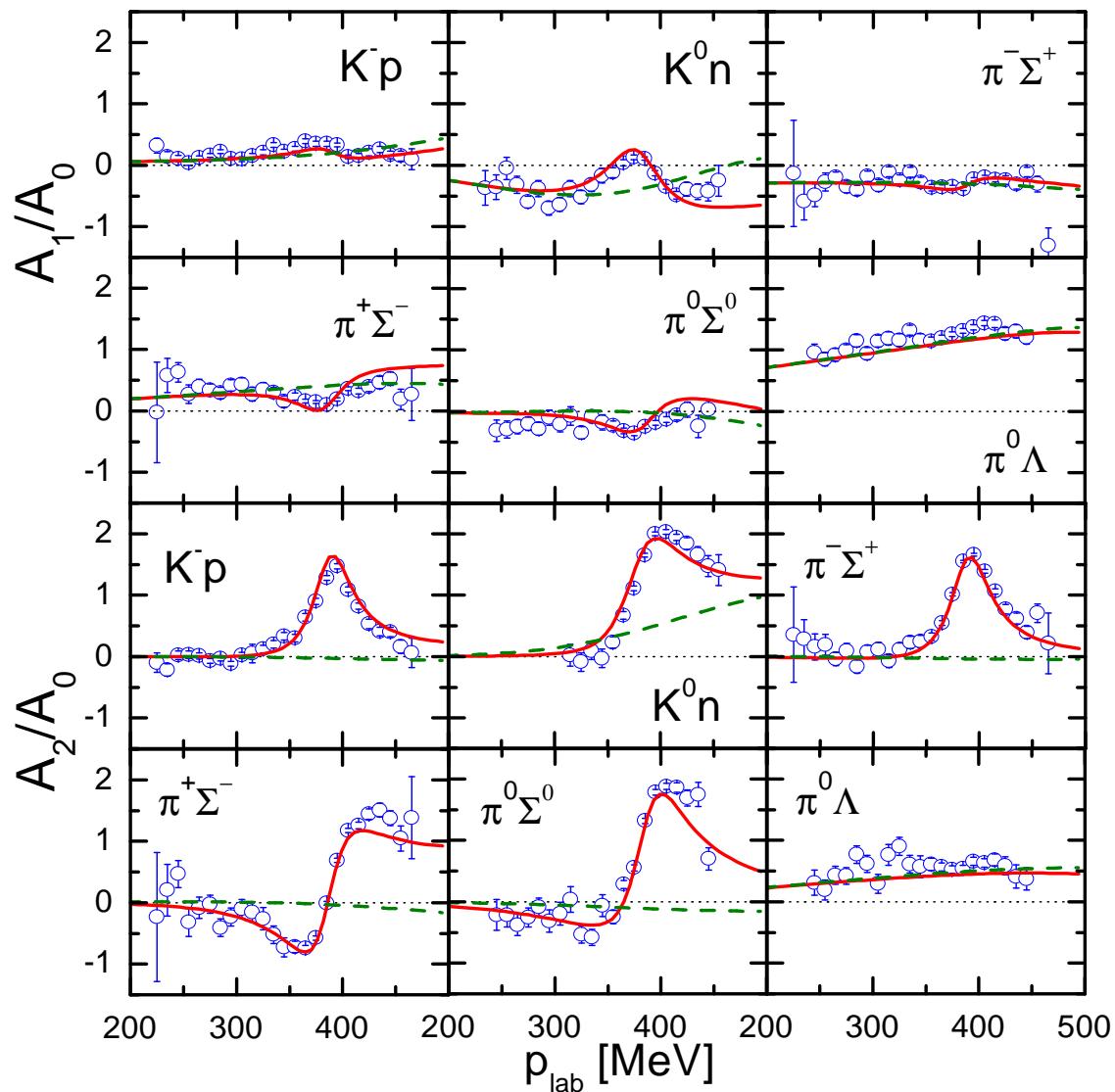
N. Kaiser, E. Oset, A. Ramos, U. Messiner . . .

M. Lutz, E. Kolomeitsev

- no  $N^*$  field is included explicitly
- V :  $SU(3)$  chiral Lagrangians up to  $(Q/\Lambda)^n$
- Numerical strategy :
  - On-shell factorization ( $N/D$  method)
  - chiral counting : u-channel  $V \rightarrow$  separable form  
 $\rightarrow$   
BS Equations  $\rightarrow$  separable  $\rightarrow$  algebraic equations
- Recent Results :
  - M. Lutz, E. Kolomeitsev :  
Can fit both  $\pi N$  and  $KN$  data ( 22 parameters)
  - T. Inoue, E. Oset et al. :  
Study  $\pi N \rightarrow \pi\pi N$  near threshold
  - . . .

Results of M. Lutz, E. Kolomeitsev :

$$\frac{d\sigma}{dcos\theta} = A_0 + A_1 P_1(\cos\theta) + A_2 P_2(\cos\theta)$$



More will be given in **M. Lutz's talk**

## Concluding Remarks

- Amplitude Analyses and Dynamical Reaction Models are complementary in  $N^*$  program
  - Has been realized in  $\Delta$  region
  - to be developed in the 2nd and 3rd  $N^*$  regions
- Gauge invariance is problematic  
Due to the need of using phenomenological form factors and/or regularization constants  
( will be discussed in Haberzettl's talk )

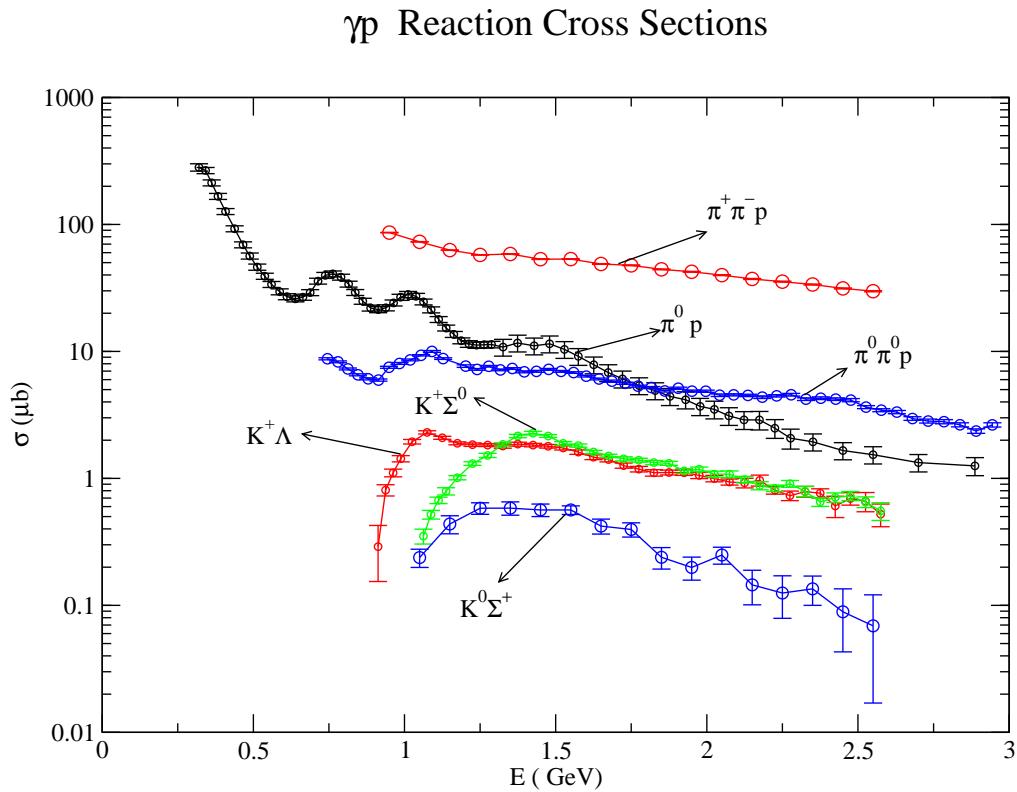
- Coupled-channel approach is **mandatory** for analyzing data of **weak** channels ( $\eta N$ ,  $KY$ ,  $\omega N$ )

Example :  $\gamma p \rightarrow K\Lambda$

**Unitarity Condition** ( $i[T - T^+] = TT^+$ )

$$Im[T_{\gamma p, K\Lambda}] = \sum_{MB} T_{\gamma p, MB} T_{K\Lambda, MB}^*$$

$$\propto \sum_{MB} \sqrt{\sigma_{\gamma p, MB}} \sqrt{\sigma_{K\Lambda, MB}}$$



$\gamma p \rightarrow (\pi N, \pi\pi N) \rightarrow K^+\Lambda$  must be significant

- collaborations are important

